

Scope of the Lecture

1. Laser amplifiers
2. Chirped pulse amplification
3. Nonlinear material response
4. Second harmonic generation
5. Optical parametric generation
6. Optical parametric oscillators

Reading: Ch. 12

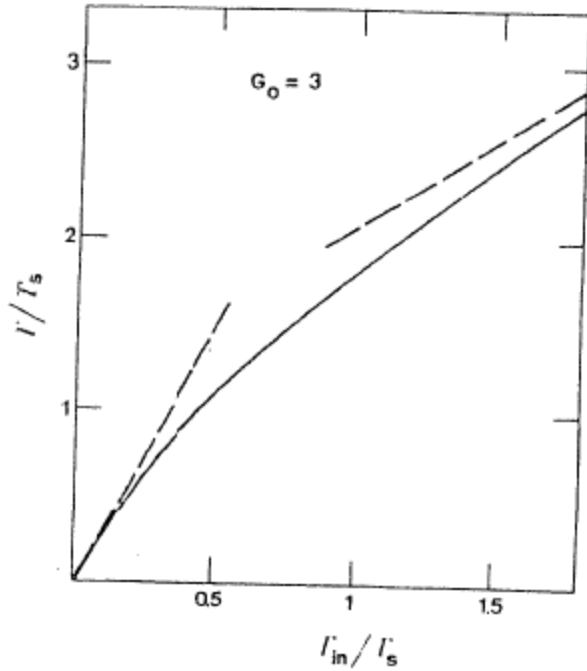
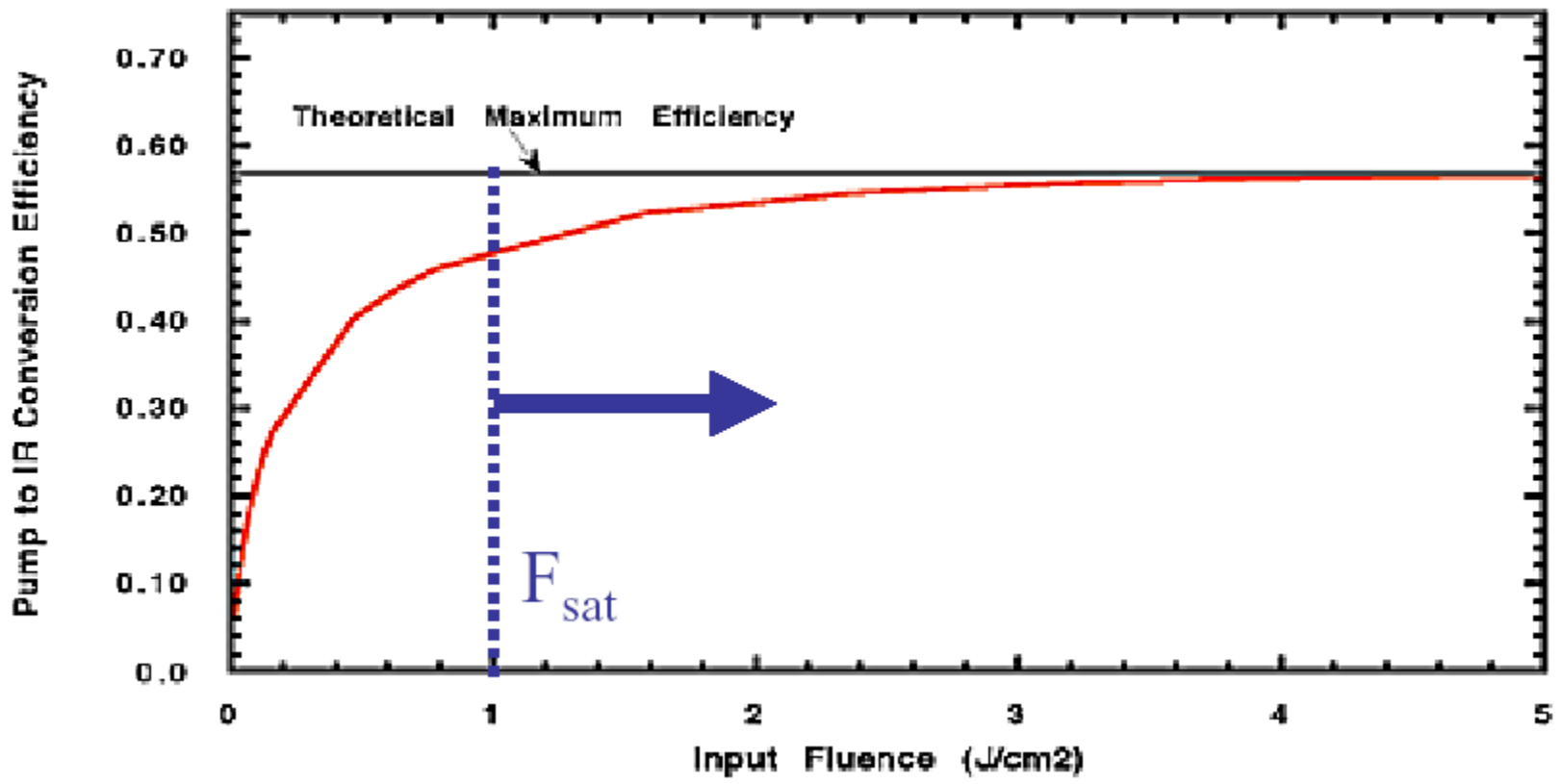
Amplifier saturation

FIG. 12.2. Output laser energy fluence Γ versus input fluence Γ_{in} for a laser amplifier with a small signal gain $G_0 = 3$. Energy fluence is normalized to the laser saturation fluence $\Gamma_s = h\nu/\sigma$.

Amplifier saturation

Operate amplifier above saturation for max efficiency



Compactness & Saturation Fluence



UCF College of Optics & Photonics/CREOL-FPCE



1-J Beam Areas at F_{sat}

RG6 & Semiconductors

$$F_{sat} \sim 1 \text{ mJ/cm}^2$$



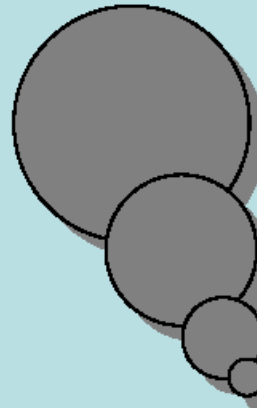
Ti:Sapphire

$$F_{sat} \sim 1 \text{ J/cm}^2$$

Yb:YAG

Yb:Glass

$$F_{sat} \sim 30\text{-}50 \text{ J/cm}^2$$



Nd:YVO4

Nd:YAG

Nd:YLF

Nd:Glass

$$F_{sat} \sim 5 \text{ J/cm}^2$$



Cr:YAG



Cr:Forsterite



Cr:LiSAF

Cr:LiCAF

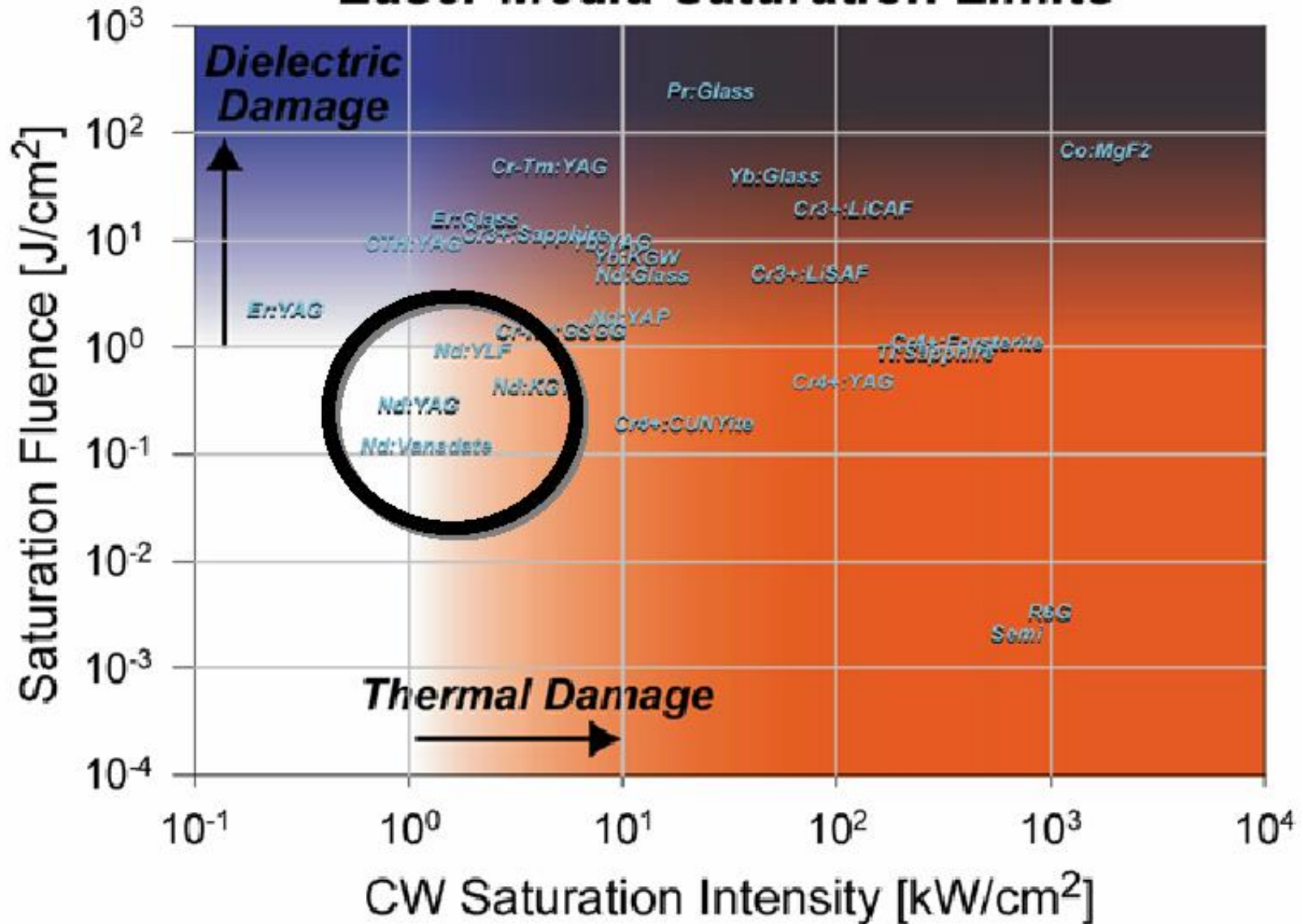


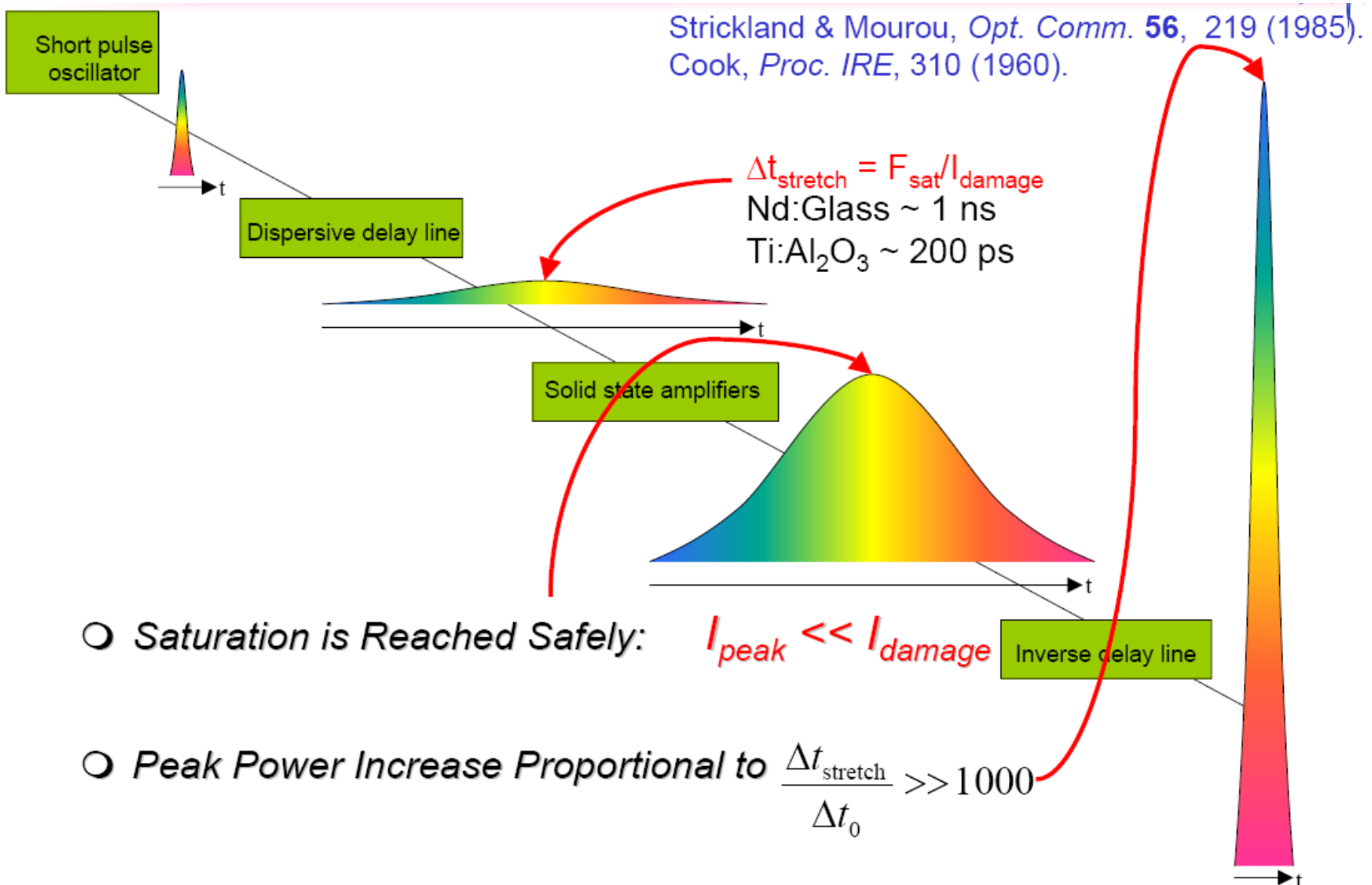
Er:YAG



Er:Glass

High F_{sat} materials
can be more compact.

High-peak versus high-average power**Laser Media Saturation Limits**

Chirped-pulse amplification (CPA)

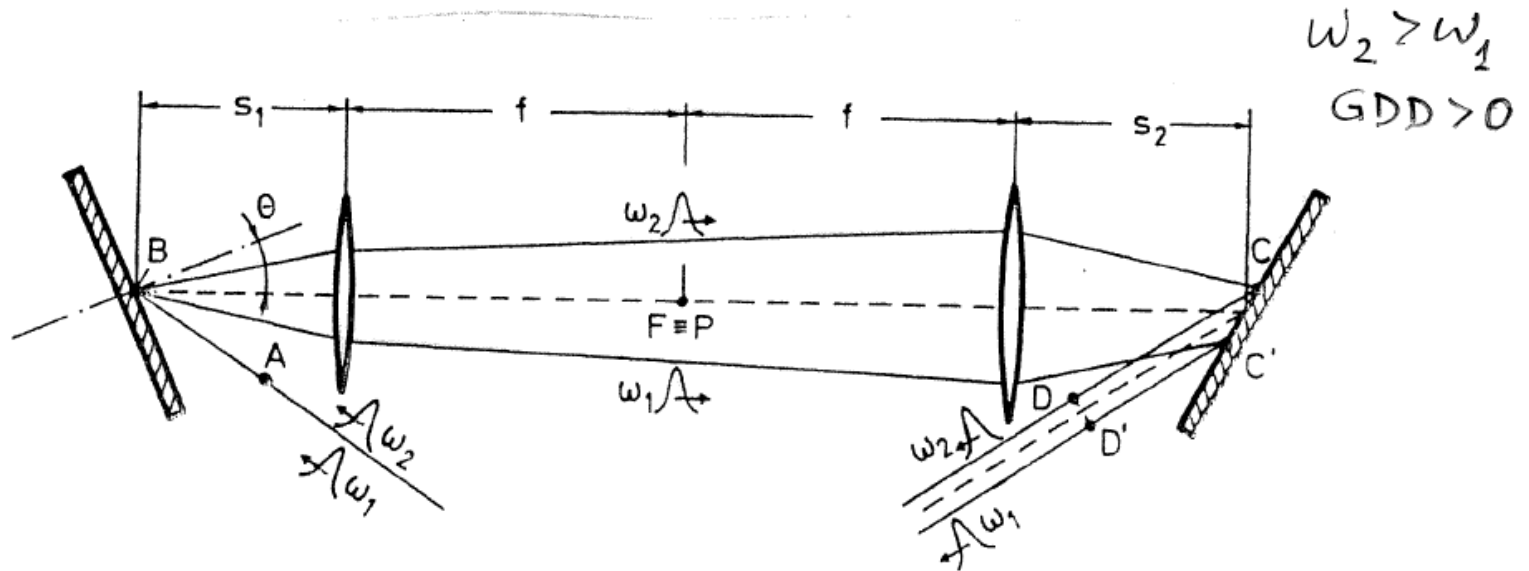
Dispersive pulse stretcher

FIG. 12.15. Pulse expander consisting of two gratings, in an antiparallel configuration, with a 1-to-1 inverting telescope between them.

Dispersive pulse compressor

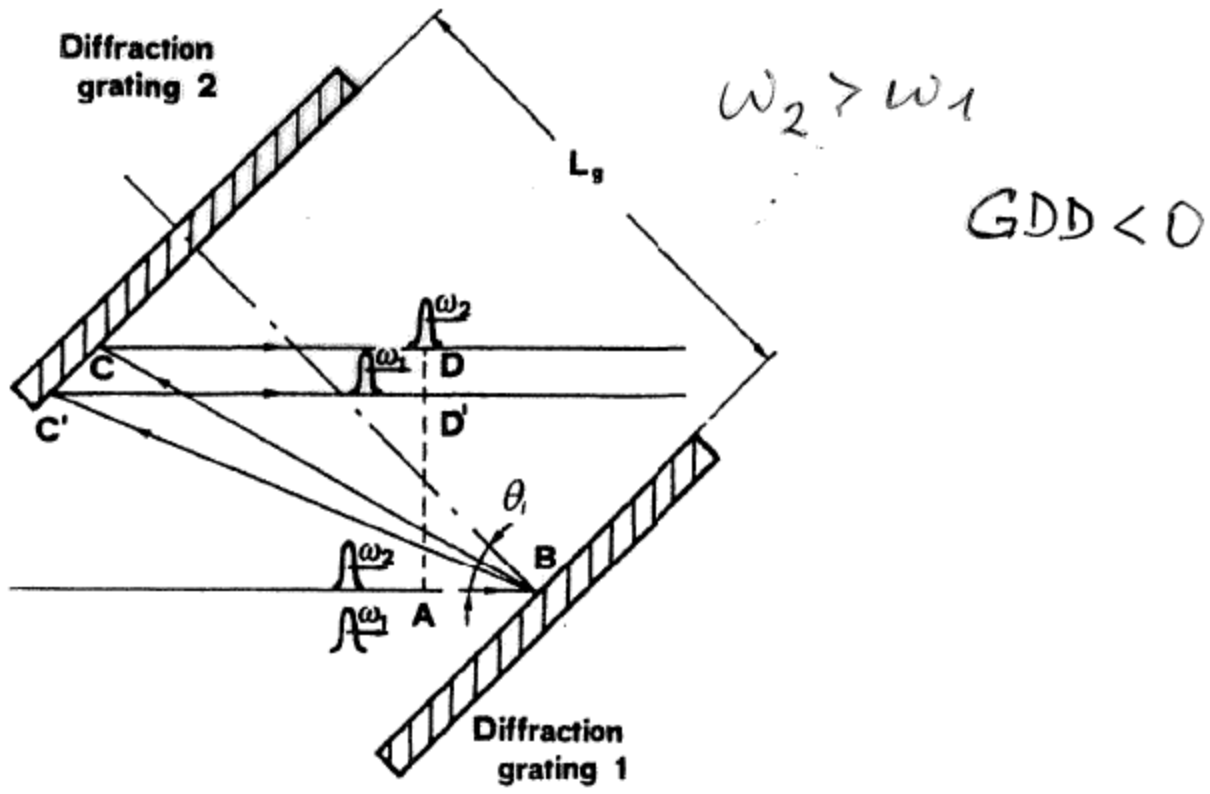
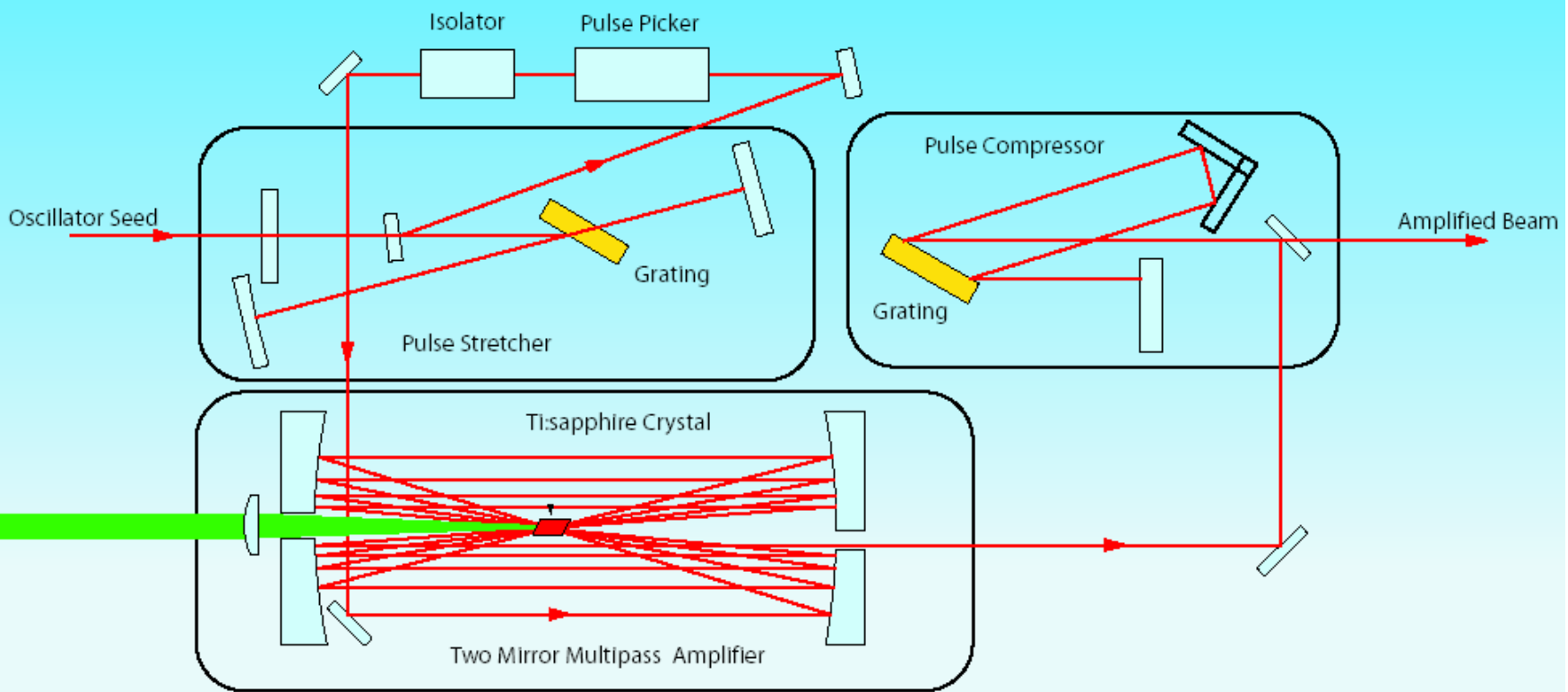
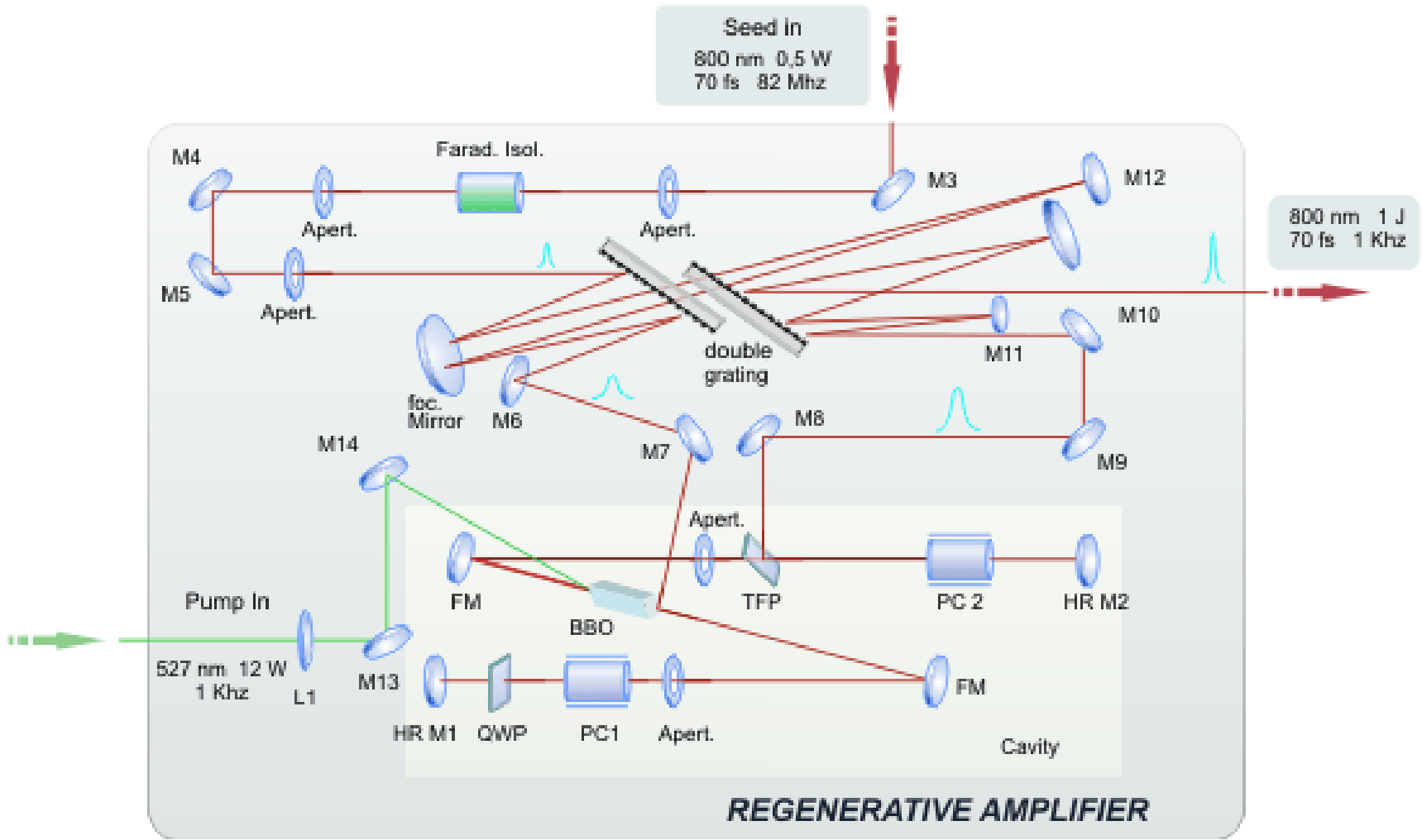


FIG. 12.14. Grating-pair for pulse compression.

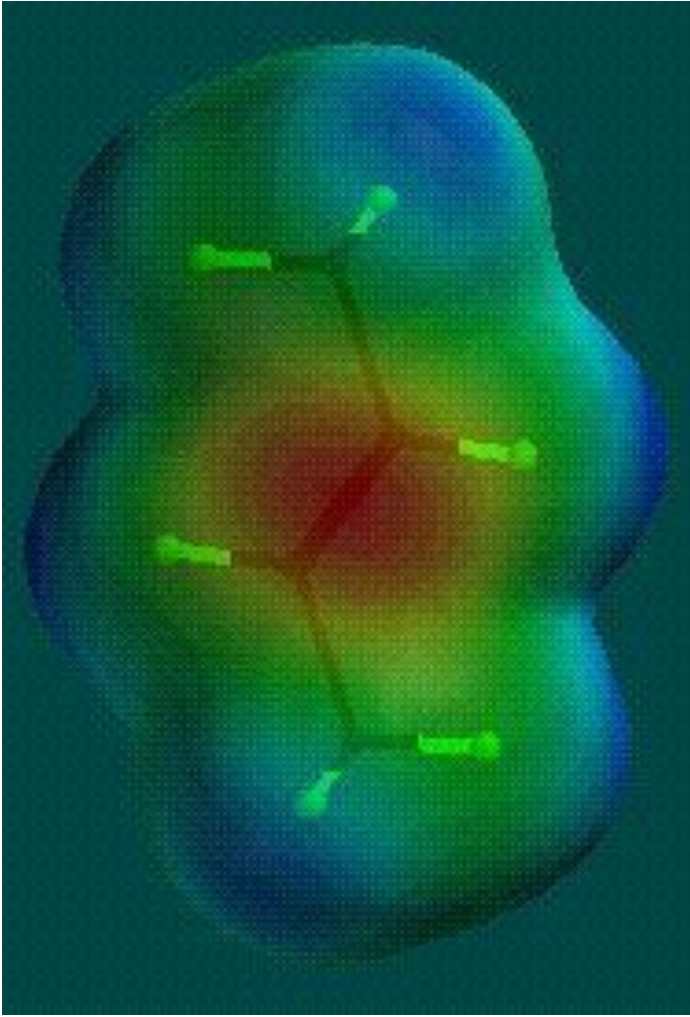
Multipass amplifier



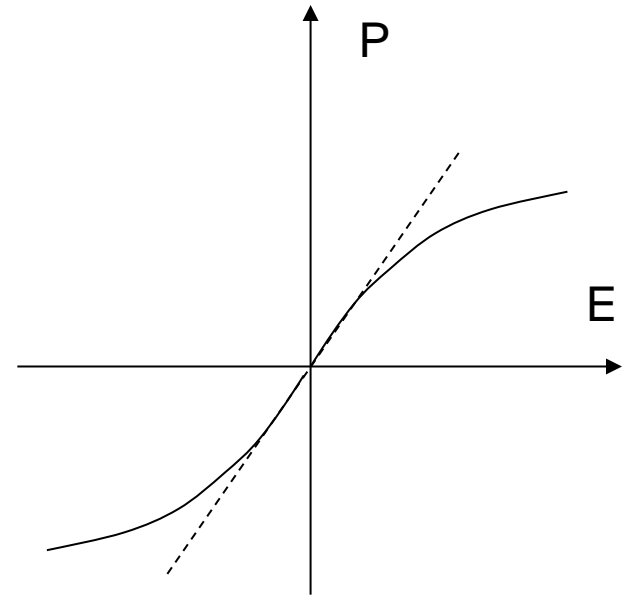
Regenerative amplifier



Nonlinear polarization



$$\vec{D} = \epsilon_0 \vec{E} + \vec{P}$$



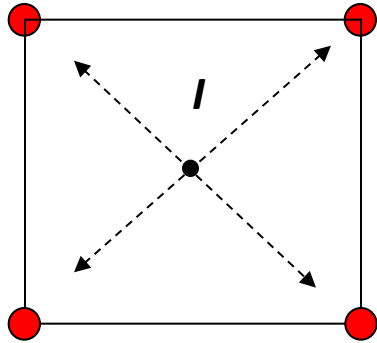
$$P = P_0 + \epsilon_0 \left(\chi^{(1)} E + \chi^{(2)} EE + \chi^{(3)} EEE + \dots \right)$$

Inversion symmetry and second-order processes

In materials with inversion symmetry 2nd order nonlinear processes are absent:

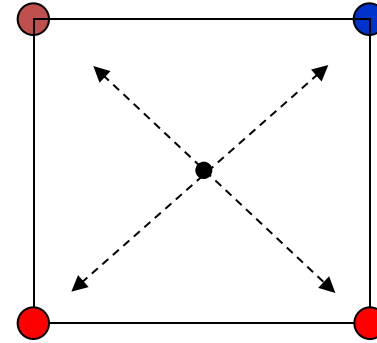
$$\chi_{ijk}^{(2)} \equiv 0$$

Inversion center I exists



Symmetry classes:
 (432).
 (622).
 (422).
 Isotropic solids.
 Atomic gasses.
 Molecular gasses.

Inversion center I absent



All remaining
 symmetry classes.
 Surfaces of
 isotropic media.

Formal proof by applying inversion operator I :

$$\mathbf{P}^{(2)} = \varepsilon_0 \chi^{(2)} : \mathbf{E}\mathbf{E}$$

$$I(\mathbf{P}^{(2)}) = \varepsilon_0 I(\chi^{(2)}) : I(\mathbf{E})I(\mathbf{E})$$

$$-\mathbf{P}^{(2)} = \varepsilon_0 I(\chi^{(2)}) : (-\mathbf{E})(-\mathbf{E}) \Rightarrow I(\chi^{(2)}) = \chi^{(2)} \Rightarrow \chi^{(2)} \equiv 0. \quad (3.4)$$

Parametric frequency conversion

Consider only second-order interactions:

$$\mathbf{P}_{NL}(\omega_1) \propto \chi^{(2)}(\omega_1; \omega_2, \omega_3) : \mathbf{E}(\omega_2) \mathbf{E}(\omega_3)$$

Processes:

- Parametric upconversion:

Sum-frequency mixing (SFM)

$$\omega_1 = \omega_2 + \omega_3, \omega_2 \neq \omega_3$$

Second harmonic generation (SHG)

$$\omega_1 = \omega_2 + \omega_3, \omega_2 = \omega_3$$

Electro-optic modulation

$$\omega_1 = \omega_2 + \omega_3, \omega_2 = 0$$

- Parametric downconversion:

Difference-frequency mixing (DFM)

$$\omega_2 = \omega_1 - \omega_3, \omega_2 \neq \omega_3$$

Parametric generation

$$\omega_2 = \omega_1 - \omega_3, \omega_2 \neq \omega_3$$

Optical rectification

$$\omega_2 = \omega_1 - \omega_3, \omega_1 = \omega_3$$

Particular process is selected by momentum conservation

$$\Delta \vec{k} = \vec{k}_1 - \vec{k}_2 - \vec{k}_3 = 0$$

$$k_j = \frac{2\pi n_j}{\lambda_j}$$

Coupled wave equations

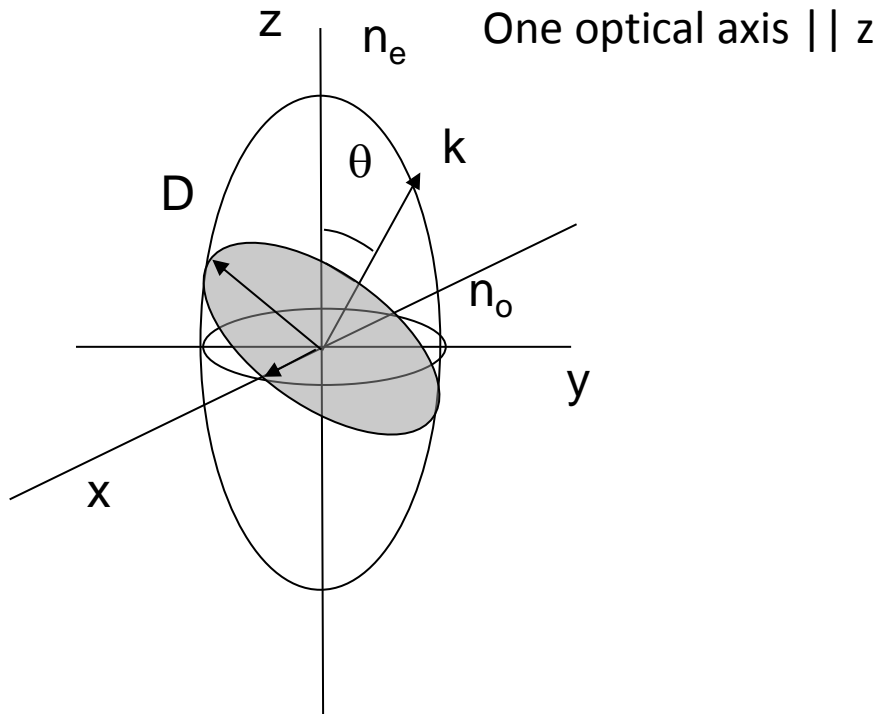
Consider [sum-frequency mixing](#) process: $\omega_1 = \omega_2 + \omega_3$; described by susceptibility $\chi_{\mu\alpha_1\alpha_2}^{(2)}(-\omega_1; \omega_2, \omega_3)$.

From Maxwell's equation, the nonlinear wave equation for ω_1 field:

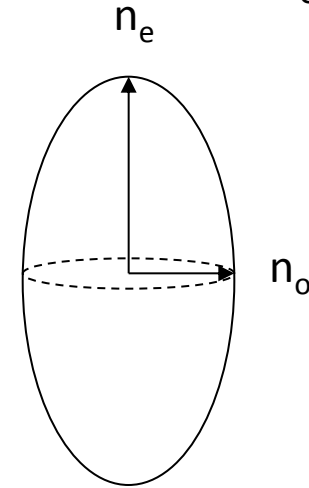
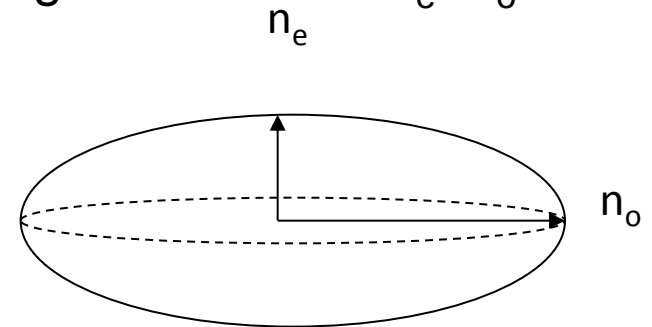
$$\frac{\partial E_1}{\partial z} \exp(i(\omega_1 t - k_1 z)) = i \frac{\omega_1 d_{\text{eff}}}{cn_1} E_2 \exp(i(\omega_2 t - k_2 z)) E_3 \exp(i(\omega_3 t - k_3 z)). \quad (3.29)$$

Need to solve coupled wave nonlinear equations for all fields:

$$\begin{aligned} \frac{\partial E_1}{\partial z} &= i \frac{\omega_1 d_{\text{eff}}}{cn_1} E_2 E_3 \exp(i\Delta k z), & \omega_1 &= \omega_2 + \omega_3; & \text{Wave-vector mismatch} \\ \frac{\partial E_2}{\partial z} &= i \frac{\omega_2 d_{\text{eff}}}{cn_2} E_1 E_3^* \exp(-i\Delta k z), & \omega_2 &= \omega_1 - \omega_3; & \Delta k = k_1 - k_2 - k_3 \\ \frac{\partial E_3}{\partial z} &= i \frac{\omega_3 d_{\text{eff}}}{cn_3} E_1 E_2^* \exp(-i\Delta k z), & \omega_3 &= \omega_1 - \omega_2; \end{aligned}$$

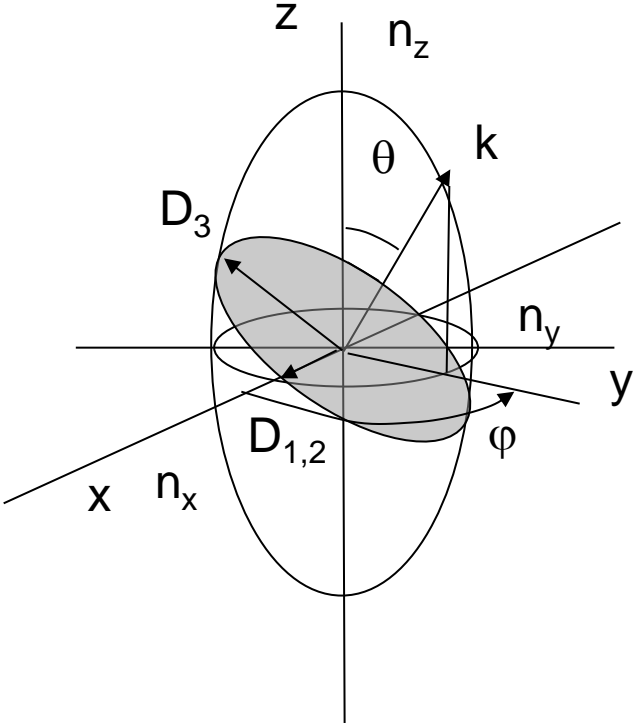
Birefringent uniaxial crystals

$$\frac{1}{[n_e(\theta, \omega_i)]^2} = \frac{\cos^2 \theta}{[n_o(0, \omega_i)]^2} + \frac{\sin^2 \theta}{[n_e(\pi/2, \omega_i)]^2}$$

Positive uniaxial $n_e > n_o$ Negative uniaxial $n_e < n_o$ 

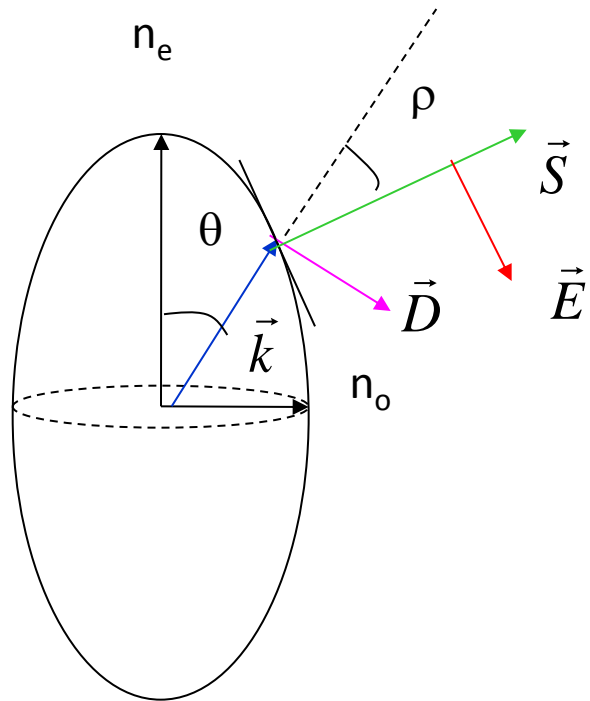
Birefringent biaxial crystals

$$n_x \neq n_y \neq n_z$$



Two optical axes

$$\frac{\sin^2 \theta \cos^2 \varphi}{n_i^{-2} - n_{x,i}^{-2}} + \frac{\sin^2 \theta \sin^2 \varphi}{n_i^{-2} - n_{y,i}^{-2}} + \frac{\cos^2 \theta}{n_i^{-2} - n_{z,i}^{-2}} = 0$$

Waves in birefringent crystals

$$\vec{D} = \epsilon_0 \vec{E} + \vec{P}$$

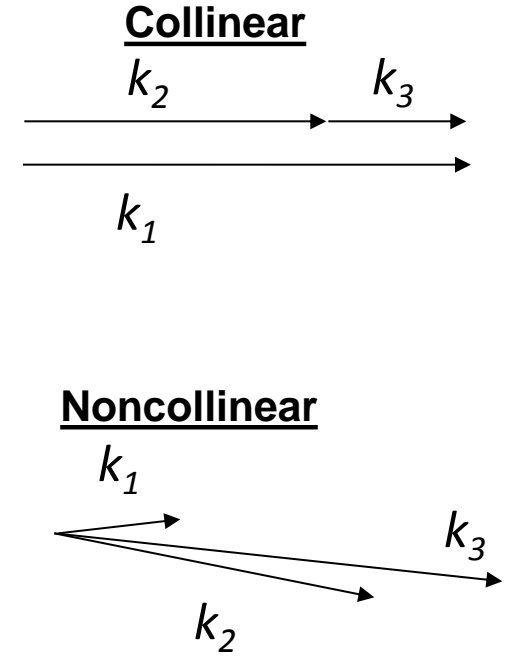
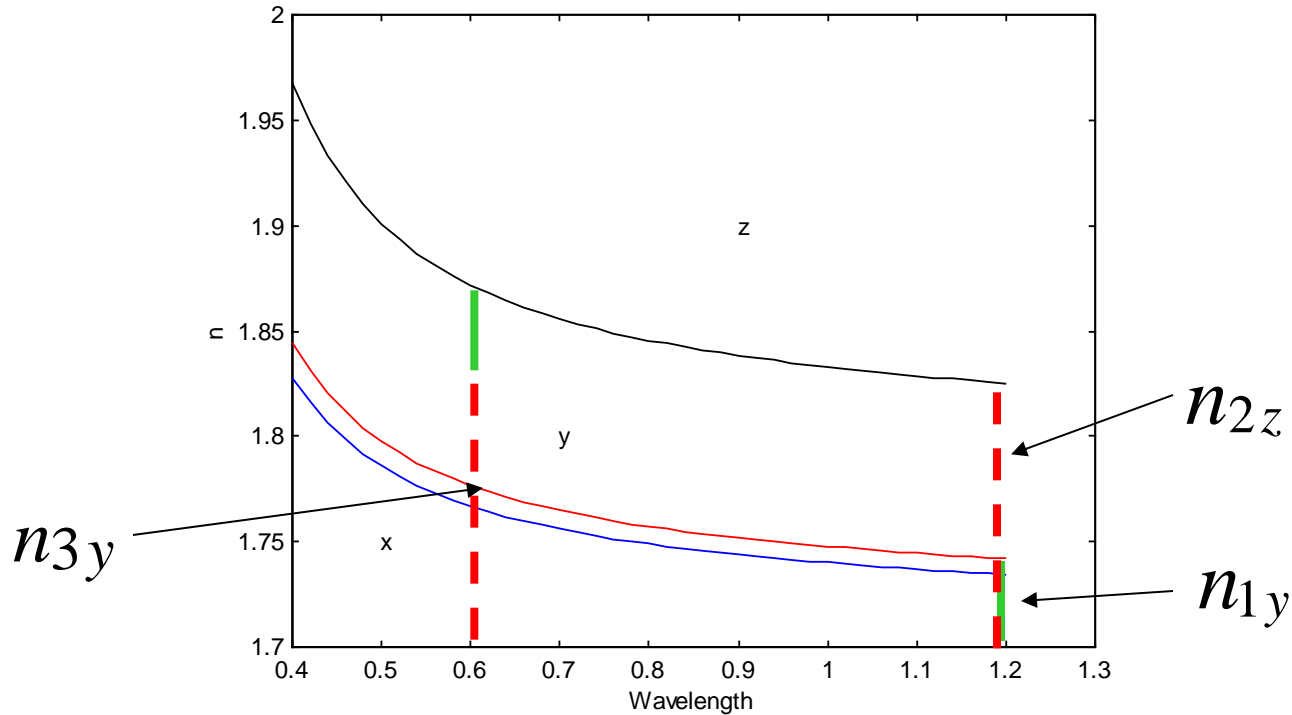
Wavevector $\vec{k} \perp \vec{D}$

Poynting vector $\vec{S} \perp \vec{E}$

Poynting vector walkoff $\rho = -\frac{1}{n} \frac{dn}{d\theta}$

Birefringence phase-matching

$$\omega_1 + \omega_2 = \omega_3, \omega_1 = \omega_2$$

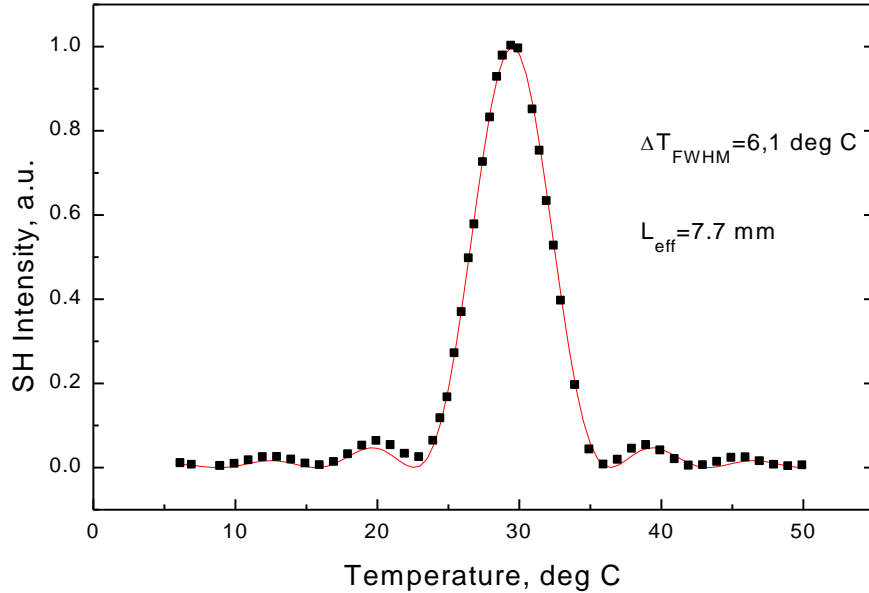


$$\frac{n_{1y}}{1.2\mu\text{m}} + \frac{n_{2z}}{1.2\mu\text{m}} = \frac{n_{3y}}{0.6\mu\text{m}} \quad \Rightarrow \quad n_{1y} + n_{2z} = 2n_{3y}$$

Momentum conservation

$$\omega_3 = \omega_1 + \omega_2$$

$$\Delta \vec{k} = \vec{k}_3 - \vec{k}_2 - \vec{k}_1 = 0$$



$$I_3(L) = \frac{\omega_3^2 d_{eff}^2}{2\epsilon_0 c^3 n_1 n_2 n_3} I_1 I_2 L^2 \text{sinc}^2(\Delta k L / 2)$$

- FWHM Phase-matching (acceptance) bandwidth: $\frac{\Delta k L}{2} \approx 0.9\pi$

- Acceptance bandwidth with respect to external parameter ψ : $\psi : \lambda, T, \theta$

$$\delta\psi = \frac{5.57}{L} \left| \frac{\partial \Delta k}{\partial \psi} \right|^{-1}$$

CWE solutions: SHG

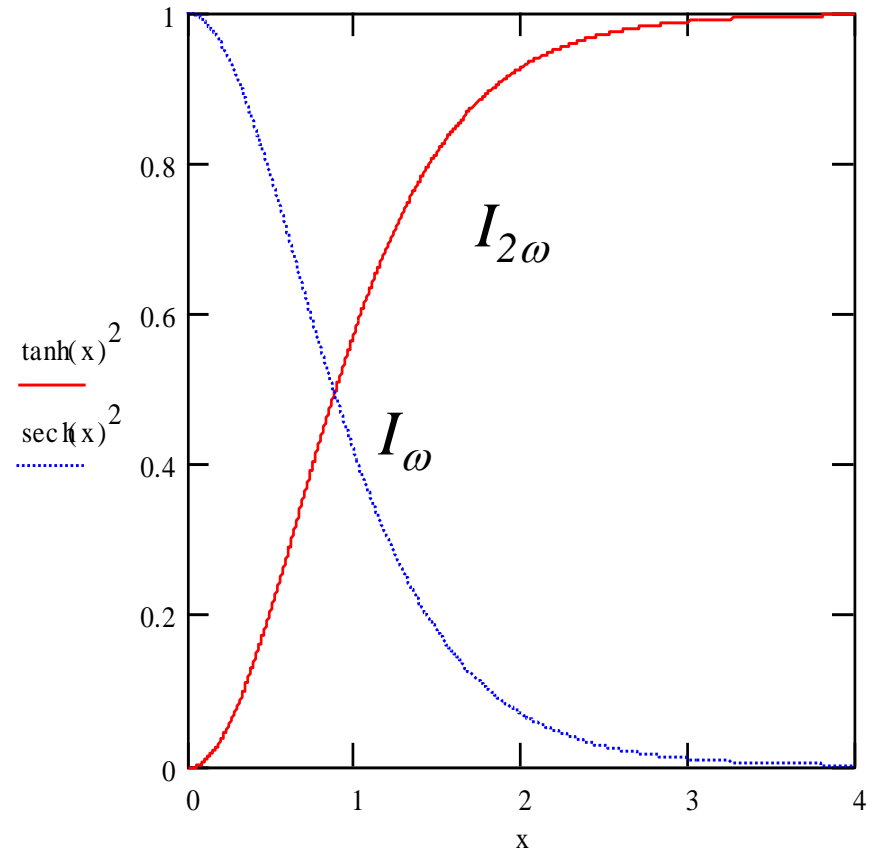
$$I_{2\omega}(L) = I_{\omega}(0) \tanh^2(\Gamma L),$$

$$I_{\omega}(L) = I_{\omega}(0) \operatorname{sech}^2(\Gamma L).$$

where
$$\Gamma = \frac{4\pi d_{eff} \sqrt{I_{\omega}(0)}}{\sqrt{2cn_{\omega}^2 n_{2\omega} \epsilon_0 \lambda_{\omega}^2}}.$$

Conversion efficiency :

$$\eta(L) = I_{2\omega}(L) / I_{\omega}(0) = \tanh^2(\Gamma L)$$



Sum-frequency mixing. Application example



Laser guide-star

Excitation of Na fluorescence with 589 nm beam focused 90 km above the Earth.

Boeing, AFRL.

CWE solutions: OPG

$$\text{Process: } \omega_i = \omega_p - \omega_s$$

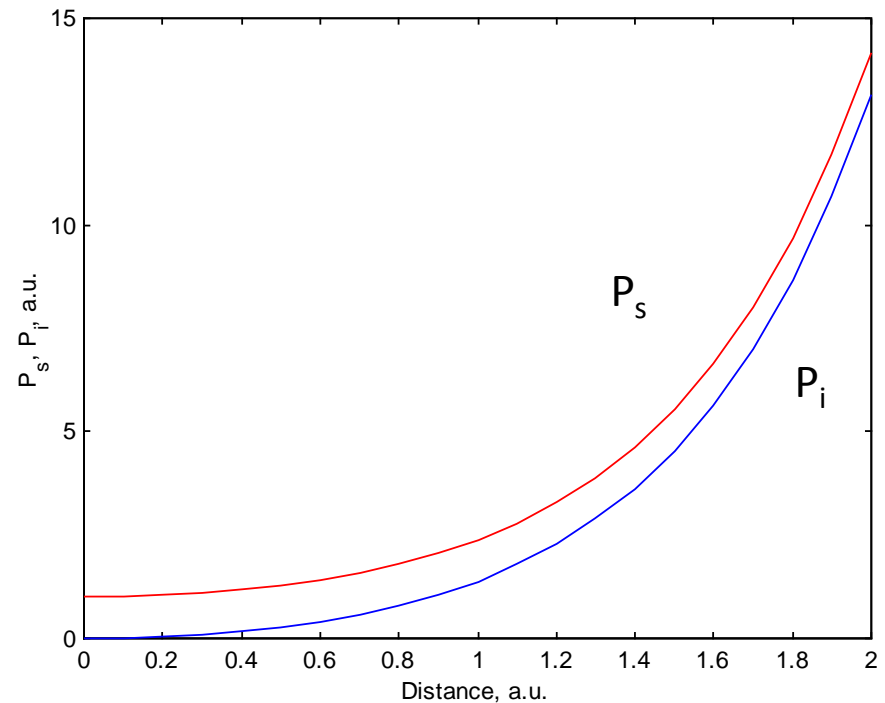
p – pump, s – signal, i - idler

$$P_s(L) / P_s(0) = \cosh^2(L / L_{NL}),$$

$$P_i(L) / P_i(0) = \sinh^2(L / L_{NL}).$$

$$\text{where: } L_{NL} = \frac{1}{4\pi d_{eff}} \sqrt{\frac{2\varepsilon_0 n_p n_s n_i c \lambda_s \lambda_i}{I_p(0)}},$$

- *Exponential growth* of signal and idler.
- Used as gain in parametric devices.

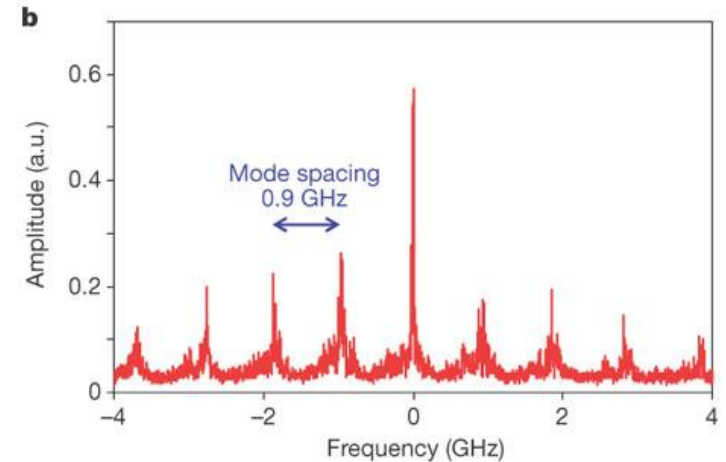
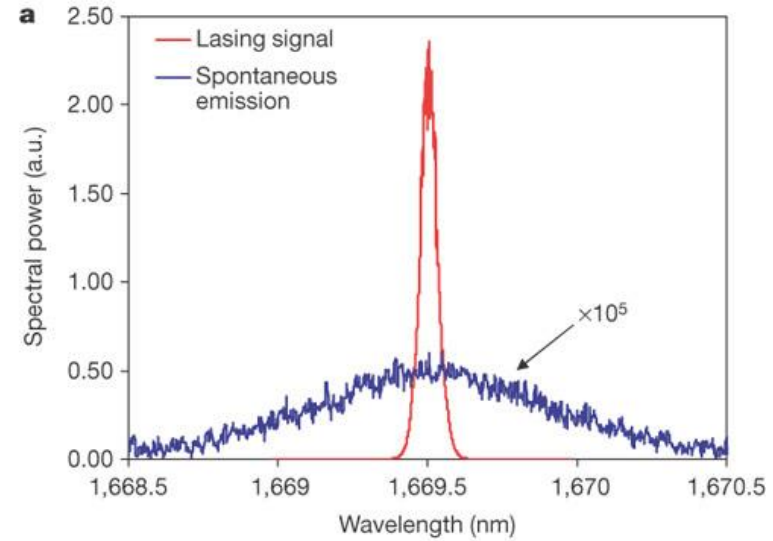
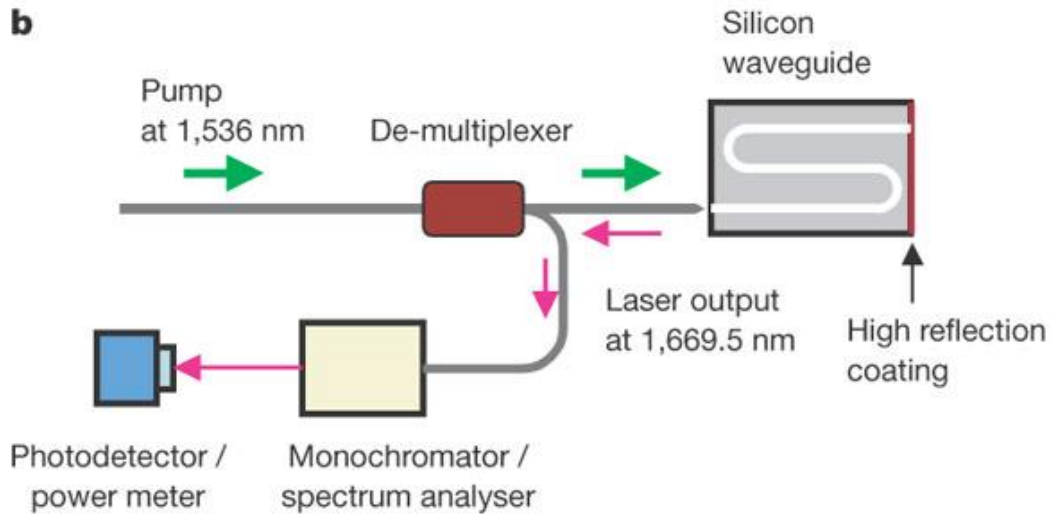
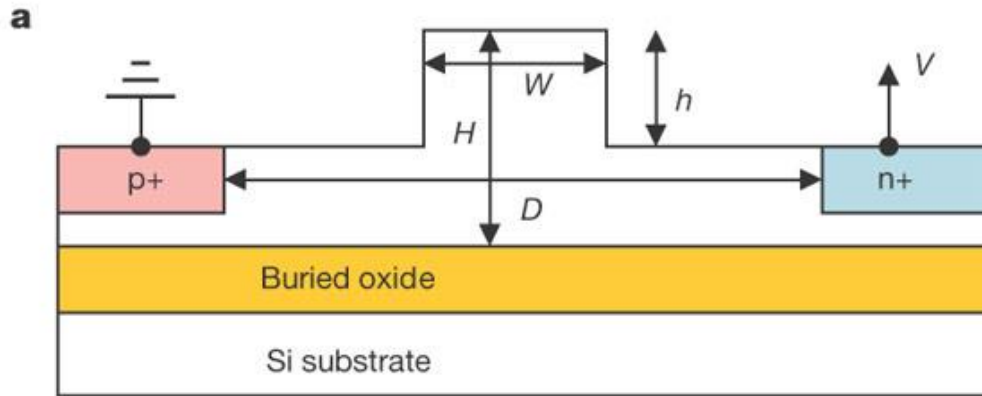


Manley-Rowe relations

$$\frac{1}{\omega_s} \frac{dP_s}{dz} = \frac{1}{\omega_i} \frac{dP_i}{dz} = -\frac{1}{\omega_p} \frac{dP_p}{dz}$$

Stimulated Raman Scattering

Recent results from Intel



Keywords

Travelling-wave amplification

Regenerative amplification

Chirped pulse amplification

Pulse stretching, compression

Nonlinear polarization

Phase-matching

Second harmonic generation

Optical parametric generation and oscillation

Problems

12.2, 12.3, 12.5, 12.7, 12.8, 12.10

Examples: 12.2, 12.3, 12.4